## Short Communication

# Small amplitude, transverse vibrations of circular plates elastically restrained against rotation with an eccentric circular perforation with a free edge 

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#### Abstract

Rayleigh-Ritz variational method is applied to the determination of the first four frequency coefficients for the title problem by making use of coordinate functions which identically satisfy the boundary conditions at the outer edge. Good stability and convergence properties are found. The mathematical model seems to be realistic and accurate, within the realm of the classical theory of vibrating plates.


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## 1. Introduction

The title problem deals with a situation which may appear in real-life vibrating systems. A small degree of eccentricity in an internal boundary may be caused by a human error in some instances. In others, it may originate by a practical reason like passage of a cable or a conduit of small diameter.

The present study makes use of the methodology employed by the senior authors (Avalos, Laura) in previous publications where the coordinate function does not satisfy the natural boundary conditions.

This is admissible since Ritz variational method is used.

## 2. Approximate analytical solution

In the case of normal modes of vibration of the vibrating system shown in Fig. 1, one takes

$$
\begin{equation*}
w^{\prime}\left(r^{\prime}, \theta, t\right)=W^{\prime}\left(r^{\prime}, \theta\right) \mathrm{e}^{\mathrm{i} \omega t} \tag{1}
\end{equation*}
$$

[^0]

Fig. 1. Vibrating mechanical system. Case of an eccentric cutout with a free edge.
Table 1
Values of the first four frequency coefficients in the case of a circular plate with a circular cutout of radius $a_{1} / a=0.1$, for different values of the non-dimensional eccentricity $e / a$ when the cutout is displaced along a radial line

| Eccentricity $e / a$ | $a /(\phi D)$ | $\Omega_{1}$ | $\Omega_{2}$ | $\Omega_{3}$ | $\Omega_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1 | 0 | 4.0053 | 13.880 | 25.416 | 29.488 |
|  | 1 | 6.0053 | 14.945 | 26.459 | 30.531 |
|  | 10 | 8.7114 | 18.513 | 30.602 | 34.989 |
|  | $\infty$ | 10.175 | 21.220 | 34.577 | 39.554 |
| 0.2 | 0 | 4.8755 | 13.855 | 25.485 | 29.383 |
|  | 1 | 6.0249 | 14.929 | 26.536 | 30.673 |
|  | 10 | 8.7407 | 18.502 | 30.706 | 35.112 |
|  | $\infty$ | 10.206 | 21.212 | 34.716 | 39.626 |
| 0.3 | 0 | 4.9030 | 13.853 | 25.573 | 29.672 |
|  | 1 | 6.0529 | 14.937 | 26.631 | 30.740 |
|  | 10 | 8.7778 | 18.534 | 30.835 | 35.165 |
|  | $\infty$ | 10.255 | 21.268 | 34.892 | 39.674 |
| 0.4 | 0 | 4.9182 | 13.871 | 25.653 | 29.636 |
|  | 1 | 6.0668 | 14.965 | 26.718 | 30.686 |
|  | 10 | 8.7909 | 18.597 | 30.956 | 35.108 |
|  | $\infty$ | 10.269 | 21.358 | 35.052 | 39.621 |
| 0.5 | 0 | 4.9249 | 13.892 | 25.706 | 29.606 |
|  | 1 | 6.0706 | 14.993 | 26.779 | 30.650 |
|  | 10 | 8.7872 | 18.651 | 31.043 | 35.108 |
|  | $\infty$ | 10.258 | 21.427 | 35.158 | 39.669 |
| 0.6 | 0 | 4.9274 | 13.904 | 25.722 | 29.622 |
|  | 1 | 6.0698 | 15.010 | 26.801 | 30.678 |
|  | 10 | 8.7744 | 18.674 | 31.078 | 35.191 |
|  | $\infty$ | 10.235 | 21.436 | 35.185 | 39.802 |
| 0.7 | 0 | 4.9272 | 13.904 | 25.700 | 29.668 |
|  | 1 | 6.0659 | 15.010 | 26.783 | 30.734 |
|  | 10 | 8.7563 | 18.653 | 31.057 | 35.279 |
| 0.8 | $\infty$ | 10.206 | 21.375 | 35.125 | 39.889 |
|  | 0 | 4.9268 | 13.8965 | 25.653 | 29.714 |
|  | 1 | 6.0620 | 14.9958 | 26.738 | 30.771 |
|  | 10 | 8.7354 | 18.5889 | 30.980 | 35.296 |
|  | $\infty$ | 10.173 | 21.2464 | 34.967 | 39.843 |

for the plate transverse displacement, and then introduces the following approximation, convenient in the case of both axisymmetric and antisymmetric modes of vibration, see for example [1].

$$
W^{\prime}\left(r^{\prime}, \theta\right) \cong W_{a}^{\prime}\left(r^{\prime}, \theta\right)=\sum_{j=0}^{j} A_{j 0}\left(\alpha_{j k} r^{\prime 4}+\beta_{j k} r^{\prime 2}+1\right) r^{\prime 2 j}
$$

$$
\begin{equation*}
+\sum_{k=1}^{K} \cos (k \theta) \sum_{j=0}^{j} A_{j k}\left(\alpha_{j k} r^{\prime 4}+\beta_{j k} r^{\prime 2}+1\right) r^{\prime j+k}, \tag{2}
\end{equation*}
$$

where $\alpha^{\prime}$ s and $\beta^{\prime}$ s of each coordinate function are determined substituting each functional relation in the governing boundary conditions at the external contour, i.e., for the elastically restrained edge against rotation,

$$
\begin{equation*}
W^{\prime}(a, \theta)=0, \quad \frac{\partial W^{\prime}\left(r^{\prime}, \theta\right)}{\partial r}(a, \theta)=\phi M_{r}\left(r^{\prime}, \theta\right), \tag{3}
\end{equation*}
$$

$M_{r}$ being the radial flexural moment, $\phi$ is the flexibility coefficient of the rotational boundary spring and $a$ is the radius of the circular plate.

The Rayleigh-Ritz variational approach requires minimization of the functional

$$
\begin{equation*}
J\left[W^{\prime}\right]=U\left[W^{\prime}\right]-T\left[W^{\prime}\right], \tag{4}
\end{equation*}
$$

where $U\left[W^{\prime}\right]$ is the maximum strain energy and $T\left[W^{\prime}\right]$ is the maximum kinetic energy for the (true) displacement amplitude $W^{\prime}$ of the plate.

Table 2
Values of the first four frequency coefficients in the case of a circular plate with a circular cutout of radius $a_{1} / a=0.2$, for different values of the non-dimensional eccentricity $e / a$ when the cutout is displaced along a radial line

| Eccentricity $e / a$ | $a /(\phi D)$ | $\Omega_{1}$ | $\Omega_{2}$ | $\Omega_{3}$ | $\Omega_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1 | 0 | 4.7199 | 13.606 | 24.881 | 30.878 |
|  | 1 | 5.9528 | 14.661 | 25.904 | 31.973 |
|  | 10 | 8.8373 | 18.152 | 29.936 | 35.440 |
|  | $\infty$ | 10.402 | 20.741 | 33.753 | 39.656 |
| 0.2 | 0 | 4.7272 | 13.673 | 25.003 | 29.899 |
|  | 1 | 5.9601 | 14.760 | 26.030 | 30.910 |
|  | 10 | 8.8355 | 18.407 | 30.112 | 35.290 |
|  | $\infty$ | 10.384 | 21.192 | 33.998 | 39.770 |
| 0.3 | 0 | 4.7467 | 13.756 | 25.311 | 29.396 |
|  | 1 | 5.9784 | 14.899 | 26.365 | 30.149 |
|  | 10 | 8.8355 | 18.738 | 30.639 | 34.365 |
|  | $\infty$ | 10.362 | 21.699 | 34.795 | 38.616 |
| 0.4 | 0 | 4.7858 | 13.833 | 25.664 | 28.983 |
|  | 1 | 6.0132 | 15.020 | 26.899 | 29.989 |
|  | 10 | 8.8483 | 18.975 | 31.481 | 34.295 |
|  | $\infty$ | 10.354 | 21.981 | 36.001 | 38.710 |
| 0.5 | 0 | 4.8273 | 13.882 | 26.141 | 27.591 |
|  | 1 | 6.0455 | 15.099 | 27.315 | 30.213 |
|  | 10 | 8.8446 | 19.072 | 31.958 | 34.861 |
| 0.6 | $\infty$ | 10.321 | 22.007 | 36.361 | 39.695 |
|  | 0 | 4.8569 | 13.895 | 26.224 | 28.910 |
|  | 1 | 6.0608 | 15.125 | 27.396 | 30.552 |
|  | 10 | 8.8068 | 19.031 | 31.905 | 35.397 |
|  | $\infty$ | 10.245 | 21.829 | 36.096 | 40.216 |
| 0.7 | 0 | 4.8749 | 13.869 | 26.125 | 29.638 |
|  | 1 | 6.0608 | 15.094 | 27.218 | 30.754 |
|  | 10 | 8.7433 | 18.858 | 31.594 | 35.529 |
|  | $\infty$ | 10.142 | 21.488 | 35.582 | 40.134 |
| 0.8 | 0 | 4.8865 | 13.830 | 25.774 | 29.629 |
|  | 1 | 6.0510 | 15.027 | 26.941 | 30.781 |
|  | 10 | 8.6676 | 18.580 | 31.151 | 35.374 |
|  | $\infty$ | 10.040 | 21.034 | 34.878 | 39.782 |

As has been shown elsewhere, see for example Ref. [2], in the case of a circular plate, each term in Eq. (4) can be written as

$$
\begin{align*}
U\left[W^{\prime}\right]= & \frac{D}{2} \iint\left\{\left[\left(\frac{\partial^{2} W^{\prime}}{\partial r^{\prime 2}}+\frac{1}{r^{\prime}} \frac{\partial W^{\prime}}{\partial r^{\prime}}\right)+\left(\frac{1}{r^{\prime 2}} \frac{\partial^{2} W^{\prime}}{\partial \theta^{2}}\right)\right]^{2}\right. \\
& -2(1-\mu)\left[\frac{\partial^{2} W^{\prime}}{\partial r^{\prime 2}}\left(\frac{1}{r^{\prime}} \frac{\partial W^{\prime}}{\partial r^{\prime}}+\frac{1}{r^{\prime 2}} \frac{\partial^{2} W^{\prime}}{\partial \theta^{2}}\right)\right. \\
& \left.\left.-\left(\frac{1}{r^{\prime}} \frac{\partial^{2} W^{\prime}}{\partial \theta \partial r^{\prime}}-\frac{1}{r^{\prime 2}} \frac{\partial W}{\partial \theta}\right)^{2}\right]\right\} r^{\prime} \mathrm{d} r^{\prime} \mathrm{d} \theta, \tag{5}
\end{align*}
$$

where $D$ is the flexural rigidity of the plate, $\mu$ its Poison's ratio and

$$
\begin{equation*}
T\left[W^{\prime}\right]=\frac{\rho \omega^{2} h}{2} \iint W^{\prime 2} r^{\prime} \mathrm{d} r^{\prime} \mathrm{d} \theta \tag{6}
\end{equation*}
$$

The integrals in expressions (5) and (6) extend over the actual area of the double connected plate under study. Introducing the non-dimensional variables

$$
\begin{equation*}
W=W^{\prime} / a, \quad r=r^{\prime} / a . \tag{7}
\end{equation*}
$$

Table 3
Values of the first four frequency coefficients in the case of a circular plate with a circular cutout of radius $a_{1} / a=0.3$ for different values of the non-dimensional eccentricity $e / a$ when the cutout is displaced along a radial line

| Eccentricity e/a | $a /(\Phi D)$ | $\Omega_{1}$ | $\Omega_{2}$ | $\Omega_{3}$ | $\Omega_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1 | 0 | 4.6631 | 12.943 | 24.116 | 34.887 |
|  | 1 | 6.1017 | 14.011 | 25.117 | 35.703 |
|  | 10 | 9.4525 | 17.521 | 29.030 | 40.993 |
|  | $\infty$ | 11.312 | 20.141 | 32.675 | 46.905 |
| 0.2 | 0 | 4.6570 | 13.293 | 24.171 | 31.376 |
|  | 1 | 6.0791 | 14.486 | 25.199 | 32.421 |
|  | 10 | 9.3097 | 18.534 | 29.260 | 37.026 |
|  | $\infty$ | 11.036 | 21.726 | 33.139 | 41.906 |
| 0.3 | 0 | 4.6503 | 13.804 | 24.502 | 29.323 |
|  | 1 | 6.0486 | 15.166 | 25.629 | 30.325 |
|  | 10 | 9.1315 | 19.795 | 30.252 | 34.611 |
|  | $\infty$ | 10.720 | 22.895 | 34.934 | 39.005 |
| 0.4 | 0 | 4.6599 | 14.236 | 26.132 | 28.533 |
|  | 1 | 6.0181 | 15.700 | 26.996 | 29.254 |
|  | 10 | 8.9600 | 20.377 | 33.413 | 33.735 |
| 0.5 | $\infty$ | 10.440 | 23.862 | 37.534 | 39.096 |
|  | 0 | 4.6601 | 14.146 | 27.519 | 28.675 |
|  | 1 | 6.0022 | 15.779 | 28.712 | 29.374 |
|  | 10 | 8.8257 | 20.066 | 32.976 | 35.728 |
|  | $\infty$ | 10.231 | 23.001 | 37.147 | 41.304 |
| 0.6 | 0 | 4.6906 | 14.143 | 27.267 | 29.322 |
|  | 1 | 6.0016 | 15.539 | 28.377 | 30.874 |
|  | 10 | 8.7232 | 19.452 | 32.618 | 36.364 |
| 0.7 | $\infty$ | 10.081 | 22.072 | 36.717 | 41.159 |
|  | 0 | 4.7296 | 13.8727 | 26.758 | 29.690 |
|  | 1 | 6.0022 | 15.2338 | 27.901 | 31.071 |
|  | 10 | 8.6237 | 18.8477 | 32.060 | 35.908 |
|  | $\infty$ | 9.9585 | 21.2702 | 35.942 | 40.352 |

Eqs. (5) and (6) above can be recast in a non-dimensional form. The functional for the whole system in Fig. 1 is

$$
\begin{align*}
\frac{2 J[W]}{D} & =\iint\left\{\left[\left(\frac{\partial^{2} W}{\partial r^{2}}+\frac{1}{r} \frac{\partial W}{\partial r}\right)+\left(\frac{1}{r^{2}} \frac{\partial^{2} W}{\partial \theta^{2}}\right)\right]^{2}\right. \\
& -2(1-\mu)\left[\frac{\partial^{2} W}{\partial r^{2}}\left(\frac{1}{r} \frac{\partial W}{\partial r}+\frac{1}{r^{2}} \frac{\partial^{2} W}{\partial \theta^{2}}\right)\right. \\
& \left.\left.-\left(\frac{1}{r} \frac{\partial^{2} W}{\partial \theta \partial r}-\frac{1}{r^{2}} \frac{\partial W}{\partial \theta}\right)^{2}\right]\right\} r \mathrm{~d} r \mathrm{~d} \theta \\
& -\Omega^{2} \iint W^{2} r \mathrm{~d} r \mathrm{~d} \theta \tag{8}
\end{align*}
$$

where as usual,
$\Omega_{i}=\sqrt{\rho h / D} \omega_{i} a^{2}$ is the non-dimensional frequency coefficient.
Minimizing the governing functional with respect to the $A_{j k}^{\prime} \mathrm{s}$ expression (8) yields a ( $J \times K$ ) homogeneous linear system of equations in the $A_{j k}^{\prime} \mathrm{s}$. A secular determinant in the natural frequency coefficients of the system results from the non-triviality condition. The present study is concerned with the determination of the first four frequency coefficients, $\Omega_{1}$ to $\Omega_{4}$ in the case of circular plates with an eccentric circular cutout.

Table 4
Values of the first four frequency coefficients in the case of a circular plate with a circular cutout of radius $a_{1} / a=0.4$ for different values of the non-dimensional eccentricity $e / a$ when the cutout is displaced along a radial line

| Eccentricity $e / a$ | $a /(\Phi D)$ | $\Omega_{1}$ | $\Omega_{2}$ | $\Omega_{3}$ | $\Omega_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1 | 0 | 4.7486 | 12.253 | 23.218 | 37.293 |
|  | 1 | 6.5515 | 13.497 | 24.234 | 38.255 |
|  | 10 | 10.689 | 17.699 | 28.203 | 42.481 |
|  | $\infty$ | 13.048 | 21.026 | 31.953 | 47.231 |
| 0.2 | 0 | 4.7058 | 12.989 | 23.507 | 35.506 |
|  | 1 | 6.4313 | 14.530 | 24.625 | 36.224 |
|  | 10 | 10.130 | 19.849 | 29.237 | 41.105 |
|  | $\infty$ | 12.085 | 24.148 | 34.248 | 42.594 |
| 0.3 | 0 | 4.64545 | 14.116 | 24.284 | 31.137 |
|  | 1 | 6.27936 | 16.008 | 25.778 | 32.196 |
|  | 10 | 9.57953 | 21.908 | 32.373 | 36.827 |
|  | $\infty$ | 11.2409 | 26.074 | 39.173 | 41.686 |
| 0.4 | 0 | 4.5813 | 15.121 | 26.056 | 29.212 |
|  | 1 | 6.1292 | 17.057 | 28.123 | 30.433 |
|  | 10 | 9.1260 | 22.259 | 33.807 | 37.166 |
|  | $\infty$ | 10.598 | 25.816 | 38.105 | 43.816 |
| 0.5 | 0 | 4.5258 | 15.310 | 27.324 | 29.620 |
|  | 1 | 6.0016 | 17.006 | 28.554 | 31.781 |
|  | 10 | 8.7818 | 21.372 | 32.890 | 38.230 |
|  | $\infty$ | 10.142 | 24.431 | 37.065 | 44.215 |
| 0.6 | 0 | 4.4953 | 14.821 | 27.367 | 30.365 |
|  | 1 | 5.9101 | 16.328 | 28.400 | 32.316 |
|  | 10 | 8.5462 | 20.144 | 32.615 | 37.799 |
| 0.7 | $\infty$ | 9.8560 | 22.835 | 36.768 | 42.757 |
|  | 0 | 4.5044 | 14.194 | 27.257 | 30.354 |
|  | 1 | 5.8716 | 15.605 | 28.266 | 31.979 |
|  | 10 | 8.4198 | 19.068 | 32.406 | 36.830 |
|  | $\infty$ | 9.7107 | 21.455 | 36.461 | 40.276 |

## 3. Numerical results

All calculations were performed for simply supported circular plates of uniform thickness, elastically restrained against rotation at the outer edge $r=a$ while the inner border of the eccentric perforation was taken to be a free edge. In all cases, the Poisson coefficient has been taken to be $\mu=0.3$.

Six tables are presented, each with a value of the ratio $a_{1} / a$ of the radius of the eccentric perforation with respect to the circular plate radius $a$. Table 1: $a_{1} / a=0.1$; Table 2: $a_{1} / a=0.2$; Table 3: $a_{1} / a=0.3$; Table 4: $a_{1} / a=0.4$; Table 5: $a_{1} / a=0.5$ and Table 6: $a_{1} / a=0.6$. In each table, in turn, four values of the nondimensional rotational spring coefficient are taken as the center of the eccentric hole is displaced along a radial line.

For the double series, Eq. (2), $J$ and $K$ up to $16 \times 6$ terms have been used, that is to say, a secular determinant of order 96 was generated for all situations. Although satisfactory convergence is already achieved for $J=8$ and $K=4$, such high values of $J$ and $K$ have been used taking advantage of the speed of modern desktop computers. As usual, special care has been taken to manipulate the numerical solving of the involved determinants, and 80 bits floating point variables (IEEE-standard temporary reals) have been used to satisfy accuracy requirements.

It is worth noting that computations are very stable and all frequency coefficients uniformly converge as the number of terms in the double series is increased.

As a general conclusion one may say that the mathematical model seems to be quite realistic and accurate, within the realm of the classical theory of vibrating plates.

Table 5
Values of the first four frequency coefficients in the case of a circular plate with a circular cutout of radius $a_{1} / a=0.5$ for different values of the non-dimensional eccentricity $e / a$ when the cutout is displaced along a radial line

| Eccentricity $e / a$ | $a /(\Phi D)$ | $\Omega_{1}$ | $\Omega_{2}$ | $\Omega_{3}$ | $\Omega_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1 | 0 | 5.0257 | 12.028 | 22.566 | 35.900 |
|  | 1 | 7.4182 | 13.802 | 23.751 | 36.921 |
|  | 10 | 12.685 | 19.979 | 28.556 | 41.463 |
|  | $\infty$ | 15.812 | 25.154 | 33.520 | 46.773 |
| 0.2 | 0 | 4.8956 | 13.161 | 23.388 | 36.826 |
|  | 1 | 7.0447 | 15.669 | 24.957 | 38.029 |
|  | 10 | 11.329 | 22.971 | 32.201 | 43.713 |
|  | $\infty$ | 13.683 | 28.003 | 40.367 | 51.604 |
| 0.3 | 0 | 4.7315 | 15.067 | 25.576 | 34.958 |
|  | 1 | 6.6535 | 17.773 | 28.369 | 36.150 |
|  | 10 | 10.265 | 24.168 | 38.129 | 41.560 |
|  | $\infty$ | 12.158 | 28.495 | 45.338 | 47.903 |
| 0.4 | 0 | 4.5685 | 16.370 | 28.769 | 31.599 |
|  | 1 | 6.3172 | 18.691 | 31.065 | 33.518 |
|  | 10 | 9.4824 | 24.044 | 36.162 | 40.835 |
|  | $\infty$ | 11.091 | 27.967 | 41.042 | 47.305 |
| 0.5 | 0 | 4.4281 | 16.491 | 27.891 | 32.311 |
|  | 1 | 6.0455 | 18.399 | 29.158 | 34.605 |
|  | 10 | 8.9203 | 23.055 | 33.606 | 40.913 |
|  | $\infty$ | 10.352 | 26.544 | 37.932 | 46.577 |
| 0.6 | 0 | 4.3183 | 15.872 | 27.113 | 32.549 |
|  | 1 | 5.8380 | 17.459 | 28.253 | 34.354 |
|  | 10 | 8.5309 | 21.556 | 32.575 | 40.053 |
| 0.7 | $\infty$ | 9.8633 | 24.591 | 36.749 | 45.532 |
|  | 0 | 4.2053 | 15.139 | 27.019 | 32.180 |
|  | 1 | 5.7245 | 16.395 | 28.072 | 33.554 |
|  | 10 | 8.2898 | 20.024 | 32.340 | 38.679 |
|  | $\infty$ | 9.5581 | 22.634 | 36.4606 | 43.430 |

Table 6
Values of the first four frequency coefficients in the case of a circular plate with a circular cutout of radius $a_{1} / a=0.6$ for different values of the non-dimensional eccentricity $e / a$ when the cutout is displaced along a radial line

| Eccentricity $e / a$ | $a /(\Phi D)$ | $\Omega_{1}$ | $\Omega_{2}$ | $\Omega_{3}$ | $\Omega_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1 | 0 | 5.5689 | 12.660 | 22.868 | 35.667 |
|  | 1 | 8.8831 | 15.709 | 24.644 | 37.019 |
|  | 10 | 15.837 | 25.503 | 32.711 | 43.372 |
|  | $\infty$ | 20.421 | 32.905 | 42.603 | 51.992 |
| 0.2 | 0 | 5.2509 | 14.693 | 24.787 | 37.480 |
|  | 1 | 7.9706 | 18.505 | 28.127 | 39.952 |
|  | 10 | 13.222 | 26.747 | 40.296 | 52.803 |
|  | $\infty$ | 16.386 | 32.547 | 49.371 | 72.388 |
| 0.3 | 0 | 4.9072 | 16.773 | 29.648 | 41.894 |
|  | 1 | 7.2083 | 19.949 | 34.096 | 43.303 |
|  | 10 | 11.441 | 26.562 | 43.330 | 49.650 |
|  | $\infty$ | 13.809 | 31.520 | 51.269 | 57.573 |
| 0.4 | 0 | 4.6021 | 17.644 | 32.721 | 35.528 |
|  | 1 | 6.6217 | 20.192 | 35.127 | 37.317 |
|  | 10 | 10.211 | 26.005 | 41.028 | 44.702 |
|  | $\infty$ | 12.102 | 30.432 | 47.097 | 52.234 |
| 0.5 | 0 | 4.2920 | 17.713 | 30.032 | 35.185 |
|  | 1 | 6.1658 | 19.760 | 31.274 | 37.328 |
|  | 10 | 9.3518 | 25.039 | 36.065 | 44.342 |
|  | $\infty$ | 10.946 | 29.018 | 40.919 | 50.946 |
| 0.6 | 0 | 4.1278 | 17.645 | 28.822 | 35.502 |
|  | 1 | 5.7953 | 18.760 | 28.910 | 36.683 |
|  | 10 | 8.7531 | 23.523 | 33.415 | 43.159 |
|  | $\infty$ | 10.160 | 27.002 | 37.735 | 49.123 |

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